

Needham, VDGf, Chapter 31, Exercise 12, p. 339.

Palmer, June 21, 2025, 9:15 am, PDT

12.

(i) *Reconsider the derivation of the Riemann tensor in [29.8], but suppose for simplicity's sake that $[\mathbf{u}, \mathbf{v}] = 0$, so that the parallelogram closes up, in which case the curvature operator simplifies to*

$$\mathbf{R}(\mathbf{u}, \mathbf{v}) = [\nabla_{\mathbf{u}}, \nabla_{\mathbf{v}}].$$

Explain why the vector holonomy of the parallelogram is given by

$$-\delta \mathbf{w}_{\parallel} = [\exp(\delta u \nabla_{\mathbf{u}}), \exp(\delta v \nabla_{\mathbf{v}})] \mathbf{w}.$$

It is given that

$$f(a + \delta x) = \exp\left[\delta x \frac{d}{dx}\right] f|_a \quad (1)$$

and by analogy that

$$\mathbf{w}(a + \mathbf{u} \delta u) = \exp[\delta u \nabla_{\mathbf{u}}] \mathbf{w}|_a \quad (2)$$

$\mathbf{w}(a + \mathbf{u} \delta u)$ in (2) is the value of $\mathbf{w}$ at the point $(a + \mathbf{u} \delta u)$, which is very near to a . This also implies that

$$\mathbf{w}(a + \mathbf{u} \delta u) \approx \delta u \nabla_{\mathbf{u}} \mathbf{w}(a) \quad (3)$$

One can deduce from lines 1 to 5 of the introduction to this exercise that $\exp(\delta u \nabla_{\mathbf{u}})$ can be rewritten as the Taylor's series:

$$\exp(\delta u \nabla_{\mathbf{u}}) \equiv \left\{ 1 + \delta u \nabla_{\mathbf{u}} + \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 + \frac{1}{3!} (\delta u)^3 (\nabla_{\mathbf{u}})^3 + \dots \right\} \quad (4)$$

and there is a parallel series for $\exp(\delta v \nabla_{\mathbf{v}})$. To help understand the operators, treat $\mathbf{w}$ as though it were $f = \{w_1, w_2, w_3\}$ and let $\nabla_{\mathbf{u}} = \left\{ \frac{d}{du_i} \right\}$. Then, by (1),

$$\mathbf{w}(a + \mathbf{u} \delta u) = \left\{ \exp\left[\delta u \frac{d}{du_i}\right] w_i \right\} |_a$$

This analysis is not needed in the proof.

The a in (2) is abstract. Letters “a”, “b”, “p”, and “o” will henceforth refer to actual points on the parallelogram such as that in [29.7]. We are also given that

$$-\delta \mathbf{w}_{||} = \delta u \delta v \mathcal{R}(\mathbf{u}, \mathbf{v}) \mathbf{w} = [\delta u \nabla_{\mathbf{u}}, \delta v \nabla_{\mathbf{v}}] \mathbf{w} = \{\delta u \nabla_{\mathbf{u}} \delta v \nabla_{\mathbf{v}} - \delta v \nabla_{\mathbf{v}} \delta u \nabla_{\mathbf{u}}\} \mathbf{w}. \quad (5)$$

This gives the difference of two instances of $\mathbf{w}_{||}$ after anticlockwise traversal on a parallelogram such as that in [29.7], first from position o *via* a to b, where we find $\mathbf{w}_{||}(b)$, then continuing on *via* p back to o, where we find $\mathbf{w}_{||}(o)$. We omit point q because the parallelogram closes in this exercise. The ordering of the two derivatives in each second derivative summand requires some derivation.

Using the $\mathbf{w}(a + \mathbf{u} \delta u)$ notation, we can show $\mathbf{w}_{||}$ on the traversal of the parallelogram.

First two sides: $\mathbf{w}_{||\text{oa}b}$ is the result of the traversal ($\mathbf{w}_{||}(o) \rightarrow \mathbf{w}_{||}(o + \mathbf{u} \delta u) \rightarrow \mathbf{w}_{||}(a + \mathbf{v} \delta v)$). The result is named $\mathbf{w}_{||}(b)$.

Second two sides: $\mathbf{w}_{||\text{bp}o}$ is the result of the traversal ($\mathbf{w}_{||}(b) \rightarrow \mathbf{w}_{||}(b - \mathbf{u} \delta u) \rightarrow \mathbf{w}_{||}(p - \mathbf{v} \delta v)$). The result is named $\mathbf{w}_{||}(o)$ to distinguish it from $\mathbf{w}(o)$.

Note that for the derivation beginning on p. 289, Needham has apparently chosen a regular vector field. That is, the fiducial vector is constant across the parallelogram, so $\mathbf{w}(o)$, $\mathbf{w}(a)$, $\mathbf{w}(b)$ and $\mathbf{w}(p)$ are all instances of the same fiducial vector anchored at different points on the traversal. Consequently, they can all be written $\mathbf{w}$, as in line 4 of the derivation. During a traversal, the parallel-transported vector rotates away from the fiducial vector field represented by $\mathbf{w}$ and is given a “||” subscript, as in $\mathbf{w}_{||}(b)$. Notice that $\delta \mathbf{w}_{||}(b)$ refers to a **difference** from the fiducial vector field at b. On the basis of the the derivation on p. 289, one can write

$$\begin{aligned} \delta \mathbf{w}_{||}(a) &\simeq \mathbf{w}_{||}(o + \mathbf{u} \delta u) - \mathbf{w} \simeq \delta u \nabla_{\mathbf{u}} \mathbf{w}(o) \\ \delta \mathbf{w}_{||}(b) &\simeq \mathbf{w}_{||}(a + \mathbf{v} \delta v) - \mathbf{w} \simeq \delta v \nabla_{\mathbf{v}} \mathbf{w}(a) \\ \delta \mathbf{w}_{||}(p) &\simeq \mathbf{w}_{||}(b - \mathbf{u} \delta u) - \mathbf{w} \simeq -\delta u \nabla_{\mathbf{u}} \mathbf{w}(p) \\ \delta \mathbf{w}_{||}(o) &\simeq \mathbf{w}_{||}(p - \mathbf{v} \delta v) - \mathbf{w} \simeq -\delta v \nabla_{\mathbf{v}} \mathbf{w}(o) \end{aligned}$$

As on p. 289, these sum to

$$\begin{aligned} -\delta \mathbf{w}_{||} &= -[\delta \mathbf{w}_{||}(a) + \delta \mathbf{w}_{||}(b) + \delta \mathbf{w}_{||}(p) + \delta \mathbf{w}_{||}(o)] \\ &\simeq \delta u \nabla_{\mathbf{u}} \mathbf{w}(o) + \delta v \nabla_{\mathbf{v}} \mathbf{w}(a) - \delta u \nabla_{\mathbf{u}} \mathbf{w}(p) - \delta v \nabla_{\mathbf{v}} \mathbf{w}(o) \\ &\simeq \delta u \delta v \{\nabla_{\mathbf{u}} \nabla_{\mathbf{v}} \mathbf{w}\} - \delta u \delta v \{\nabla_{\mathbf{v}} \nabla_{\mathbf{u}} \mathbf{w}\} \text{ (by last line, p. 289; intermediate step omitted)} \\ &= \{\delta u \nabla_{\mathbf{u}}\} \circ \{\delta v \nabla_{\mathbf{v}}\} \mathbf{w} - \{\delta v \nabla_{\mathbf{v}}\} \circ \{\delta u \nabla_{\mathbf{u}}\} \mathbf{w} \\ &\simeq \{\exp[\delta u \nabla_{\mathbf{u}}] \times \exp[\delta v \nabla_{\mathbf{v}}] - \exp[\delta v \nabla_{\mathbf{v}}] \times \exp[\delta u \nabla_{\mathbf{u}}]\} \mathbf{w} \end{aligned}$$

[By (2), (3), and (4) above, substituting products of exponential operator for $\{\delta u \nabla_{\mathbf{u}}\} \circ \{\delta v \nabla_{\mathbf{v}}\}$.]

Because derivatives factorize and obey the Index Law $D^m . D^n u = D^{m+n} . u$ (Paggio, H.T.H., *Elementary Treatise...*, 1949, p. 31), they must also follow the rules of the Riemann bracket notation. That is, one can write $D^m . E^n u$ and $[D, E] u = D . Eu - E . Du$, where D and E are derivatives, such as $\frac{d}{dx}$ or $\nabla_{\mathbf{u}}$. The exponential operator is actually a derivative, or more precisely, a sum of derivatives, because it is a Taylor's series applied to the gradient. Then one can write

$$-\delta \mathbf{w}_{\parallel} = [\exp(\delta u \nabla_{\mathbf{u}}), \exp(\delta v \nabla_{\mathbf{v}})] \mathbf{w}$$

This is what we wanted.

(ii) Deduce that

$$-\delta \mathbf{w}_{\parallel} = \delta u \delta v \mathcal{R}(\mathbf{u}, \mathbf{v}) \mathbf{w} + (\text{third-order error})$$

where the “third-order error” is made up of terms involving $(\delta u)^p (\delta v)^q$, where $(p + q) \geq 3$.

In part (i) we found that

$$-\delta \mathbf{w} \simeq \exp(\delta u \nabla_{\mathbf{u}}) \cdot \exp(\delta v \nabla_{\mathbf{v}}) \mathbf{w} - \exp(\delta v \nabla_{\mathbf{v}}) \cdot \exp(\delta u \nabla_{\mathbf{u}}) \mathbf{w} \quad (5)$$

and

$$\exp(\delta u \nabla_{\mathbf{u}}) \equiv \left\{ 1 + \delta u \nabla_{\mathbf{u}} + \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 + \frac{1}{3!} (\delta u)^3 (\nabla_{\mathbf{u}})^3 + \dots \right\}$$

and there is a parallel series for $\exp(\delta v \nabla_{\mathbf{v}})$. Then by composing the terms in the series, the first term $\exp(\delta u \nabla_{\mathbf{u}}) \cdot \exp(\delta v \nabla_{\mathbf{v}})$ in (5) becomes

$$\begin{aligned} & 1 \cdot \left\{ 1 + \delta v \nabla_{\mathbf{v}} + \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} \\ & + \delta u \nabla_{\mathbf{u}} \circ \left\{ 1 + \delta v \nabla_{\mathbf{v}} + \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} \\ & + \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 \circ \left\{ 1 + \delta v \nabla_{\mathbf{v}} + \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} \\ & \dots \\ & = \left\{ 1 + \delta v \nabla_{\mathbf{v}} + \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} + \left\{ \delta u \nabla_{\mathbf{u}} + \delta u \nabla_{\mathbf{u}} \circ \delta v \nabla_{\mathbf{v}} + \delta u \nabla_{\mathbf{u}} \circ \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} \\ & + \left\{ \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 + \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 \circ \delta v \nabla_{\mathbf{v}} + \frac{1}{2!} (\delta u)^2 (\nabla_{\mathbf{u}})^2 \circ \frac{1}{2!} (\delta v)^2 (\nabla_{\mathbf{v}})^2 + \dots \right\} \end{aligned}$$

Write just the terms where $p + q \leq 2$ and subtract the second term $\exp(\delta v \nabla_{\mathbf{v}}) \cdot \exp(\delta u \nabla_{\mathbf{u}})$ in (3).

$$\begin{aligned}
&= \{1 + \delta v_{\nabla_v} + \delta u_{\nabla_u} + \delta u_{\nabla_u} \circ \delta v_{\nabla_v} + (\text{terms with } p + q \geq 3)\} \\
&- \{1 + \delta v_{\nabla_v} + \delta u_{\nabla_u} + \delta v_{\nabla_v} \circ \delta u_{\nabla_u} + (\text{terms with } p + q \geq 3)\} \\
&= \{\delta u_{\nabla_u} \circ \delta v_{\nabla_v} - \delta v_{\nabla_v} \circ \delta u_{\nabla_u} + (\text{terms with } p + q \geq 3)\}
\end{aligned}$$

because in general, $\delta u_{\nabla_u} \circ \delta v_{\nabla_v} \neq \delta v_{\nabla_v} \circ \delta u_{\nabla_u}$.

This shows that

$$\begin{aligned}
&\{\exp(\delta u_{\nabla_u}) \cdot \exp(\delta v_{\nabla_v}) - \exp(\delta v_{\nabla_v}) \cdot \exp(\delta u_{\nabla_u})\} \mathbf{w} \\
&= \{\delta u_{\nabla_u} \circ \delta v_{\nabla_v} - \delta v_{\nabla_v} \circ \delta u_{\nabla_u} + (\text{terms with } p + q \geq 3)\} \mathbf{w} \\
&= \{[\delta u_{\nabla_u}, \delta v_{\nabla_v}] + (\text{terms with } p + q \geq 3)\} \mathbf{w}
\end{aligned}$$

From (29.5) and the first box on p. 288, one can write:

$$[\delta u_{\nabla_u}, \delta v_{\nabla_v}] \mathbf{w} = (\delta u \delta v) \mathcal{R}(\mathbf{u}, \mathbf{v}) \mathbf{w}$$

It follows that

$$-\delta \mathbf{w} = (\delta u \delta v) \mathcal{R}(\mathbf{u}, \mathbf{v}) \mathbf{w} + (\text{third - order error})$$